

PHYS 393  
Low Temperature Physics  
Set 4:

Superconductivity – Theory

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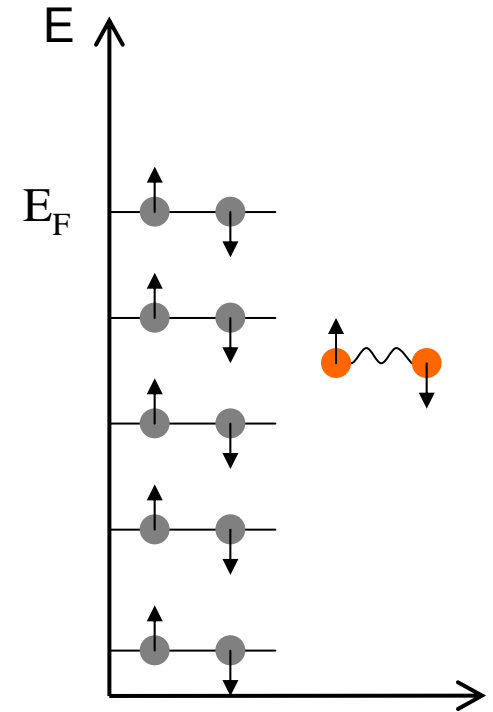
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# Superconductivity theory

- Many efforts over decades resulted in models with limited success.
- Breakthrough in 1957 (Bardeen, Cooper, Schrieffer), **BCS theory**:
  - Pairs of electrons (fermions) can couple to have zero or integer angular momentum and thus become bosons (Cooper pairs, 1956)
  - Pairing interaction mechanism (electron-phonon-electron) proposed in 1950 independently by Fröhlich (Liverpool) and Bardeen

# Cooper pair formation

- Fermi sea: all states with  $E < E_F$  (Fermi energy) occupied
- Addition of non-interacting electron pair will stay on surface ( $E_F$ )
- In the presence of small attractive interaction between pair they will "immerse" in Fermi sea as the attractive interaction lowers the pair's energy
- **Question:** what can create attractive interaction between two electrons (overcoming the strong Coulomb repulsion)?



Sketch of Fermi gas (electrons inside conductor) and a Cooper pair (right)

# Pairing interaction

- Electron flying through lattice of positive ions exerts attractive force on them creating a positive cloud in its wake
- This surplus of positive charge attracts another electron
- Development of ion cloud delayed by time of order of lattice vibration period ( $10^{-13}\text{s}$ )
- Electrons move typically with  $10^8\text{cm/s}$
- Thus first electron has moved by  $100\text{nm}$  before second electron is attracted; distance reduces Coulomb repulsion leaving net attractive force
- Electron-cloud-electron moves coupled through the material: Cooper pair
- Interaction described as exchange of virtual phonons

# Cooper pairs

- Phonon exchange most efficient for vanishing center-of-mass momentum of the pair.
- Hence electrons have wavevectors  $\underline{k}$  and  $-\underline{k}$
- This corresponds to space wavefunction with zero angular momentum:  $L=0$
- Spin wavefunction is antisymmetric ( $\uparrow\downarrow$ ):  $S=0$
- Overall wavefunction antisymmetric under particle exchange (made of fermions)
- But total angular momentum zero:  $\underline{J}=\underline{L}+\underline{S}=0+0=0$
- Hence pair acts as a boson

# BCS Ground State

- Many bosons can occupy the ground state
- Correlated wavefunction of pairs: **single wavefunction for all as in Bose-Einstein condensate**
- Identify superconducting current with that carried by Cooper pairs
- Normal current carried by single electrons in excited states
- Ground state separated from first excited state by energy gap  $\Delta$
- First excited state: broken pairs: for each pair we get two electron-hole pairs

# Summary points

- Superconductivity depends on condensation of bosons (Cooper pairs) into ground state
- Critical temperature  $T_C$  set by strength of pairing interaction, not by B-E condensation ( $T_B \gg T_C$ )
- Magnetic field destroys superconductivity as its action is to align electron spins, breaking up  $\uparrow\downarrow$  pairs
- BCS theory relates energy gap at  $T=0$   $\Delta(0)$  to electron-lattice interaction strength  $V$  via:

$$\Delta(0) = 1.76 k T_C = 2 \hbar \omega_D e^{-\frac{1}{g(\mu)V}}$$

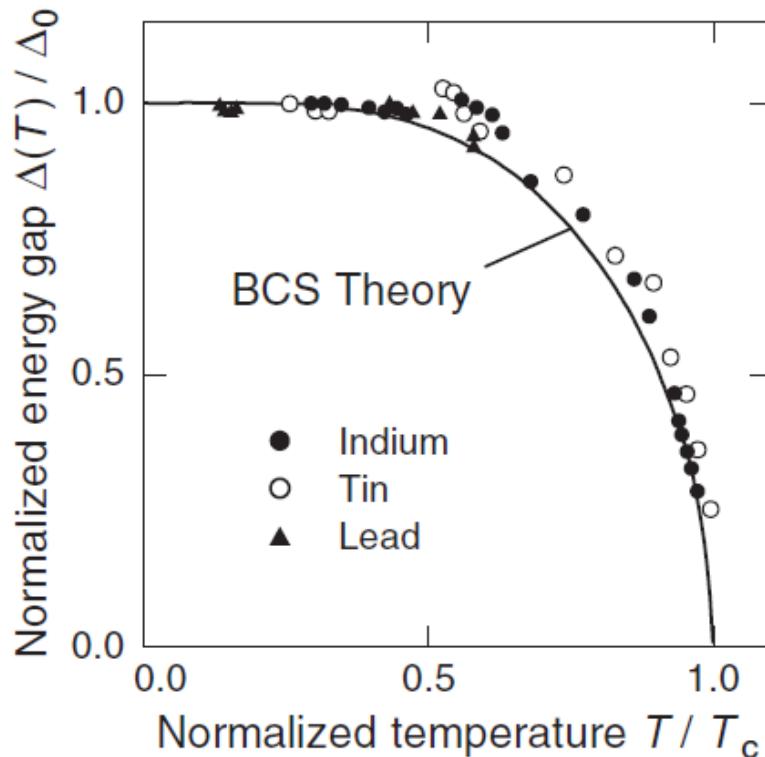
Density of states  
at Fermi surface

Interaction  
strength

# Comments

- For metals the model factor 1.76 is in the range 1.6-2.15: reasonable agreement
- Energy gap  $\Delta$  effectively constant in region  $0 < T < T_c/2$
- For  $T_c/2 < T < T_c$  it varies as:

$$\Delta(T) \approx 3.1kT_c \left[ 1 - \frac{T}{T_c} \right]^{1/2}$$





# On interaction strength

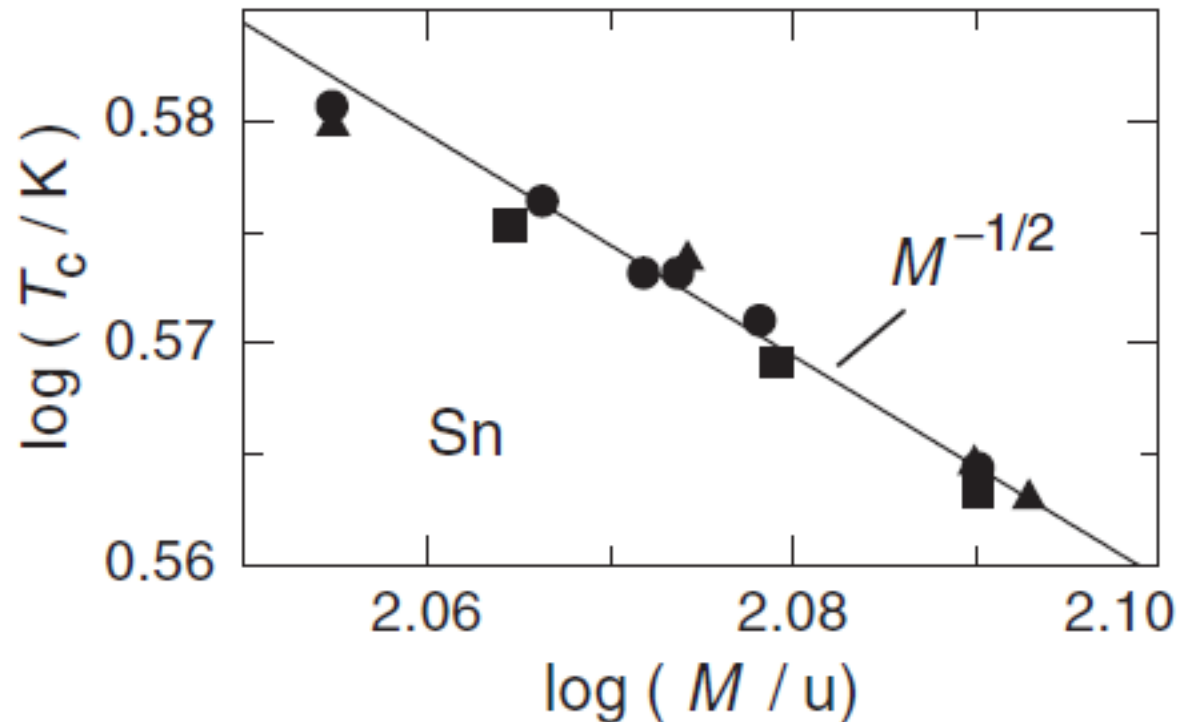
$$\Delta(0) = 1.76 kT_c = 2\hbar\omega_D e^{-\frac{1}{g(\mu)V}}$$

- Large electron-lattice interaction  $V$  related to high  $T_c$
- In normal state large  $V$  means large resistivity
- This is the connection between superconductivity and poor metallic conductivity in normal state
- This is why good conductors do not show superconductivity

# Isotope effect

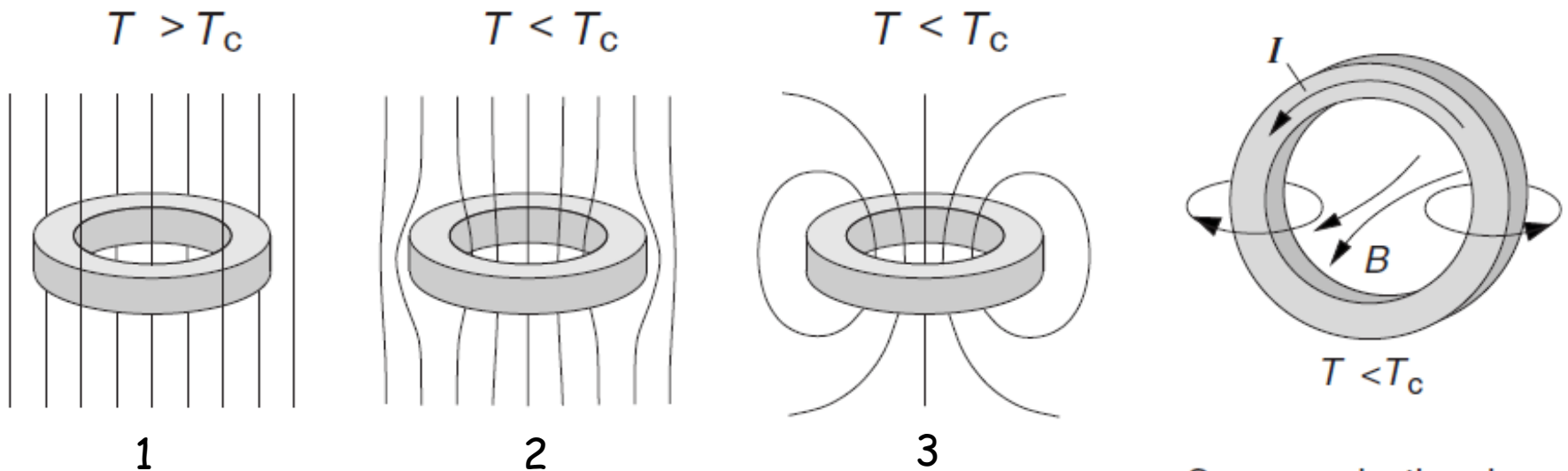
- First evidence that lattice is involved in pairing interaction
- Measurements of  $T_c$  for different isotopes of same element versus isotope mass  $M$ :

$$T_c \propto \sqrt{\frac{1}{M}}$$



# Persistent currents

1. Magnetic field at  $T > T_c$  (normal conducting state) penetrates material
  2. Temperature lowered below  $T_c$ : magnetic field expelled from material (Meissner effect)
  3. Magnetic field switched off: magnetic flux cannot be expelled from interior of ring as it would have to cross material: magnetic flux remains frozen
- Frozen flux  $\leftrightarrow$  persistent current



Superconducting ring  
with persistent current

# Coherence and macroscopic wavefunction

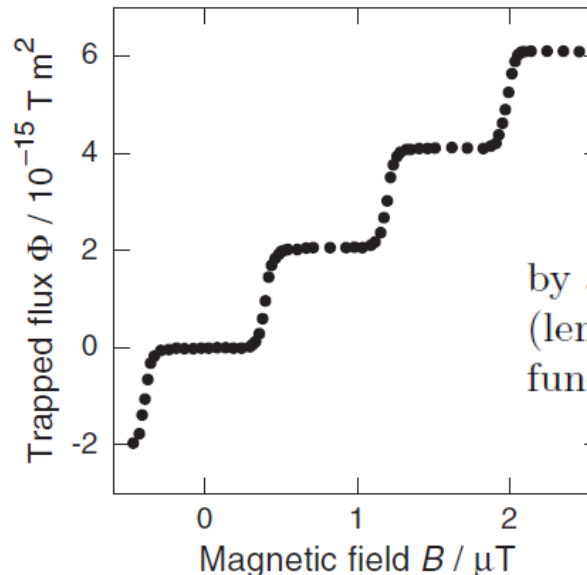
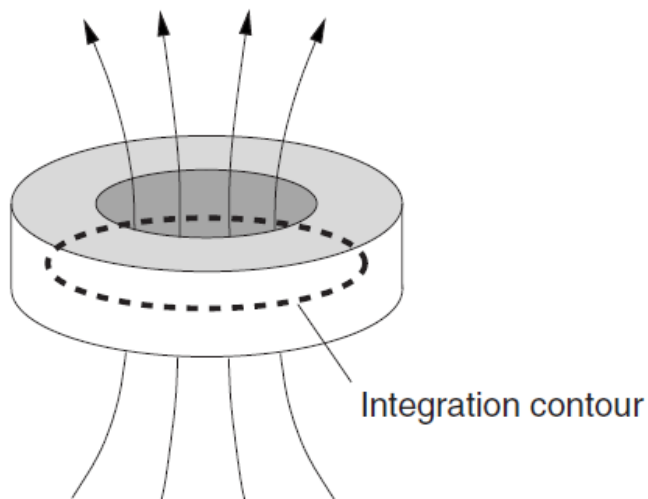
- Maximum energy uncertainty for Cooper pairs is  $2\Delta$ , leading to position uncertainty of order **1000nm (coherence length)**
- Large than Cooper pair size
- Leading to large space overlaps
- Implying high correlation between pairs
- Indeed in B.E.C. (ground state) all Cooper pairs belong to **common quantum state**
- This is described by **macroscopic wavefunction**  
 **$\Psi = \Psi_0 \exp(i\phi(r))$**
- There is a well-defined phase in the whole sample

# Flux quantization (I)

- Macroscopic wavefunction has unique phase
- Wavefunction is single-valued
- Phase change for a **closed contour** must be  $2\pi n$
- Leading (after a bit of algebra) to:

$$\Phi = n\Phi_0 \quad \text{with flux quantum} \quad \Phi_0 = \frac{h}{2e}$$

In agreement with charge of current carrier (Cooper pair)



Magnetic flux trapped by a thin hollow cylinder made of tin (length 24 mm, diameter 56  $\mu\text{m}$ ) as a function of the cooling field

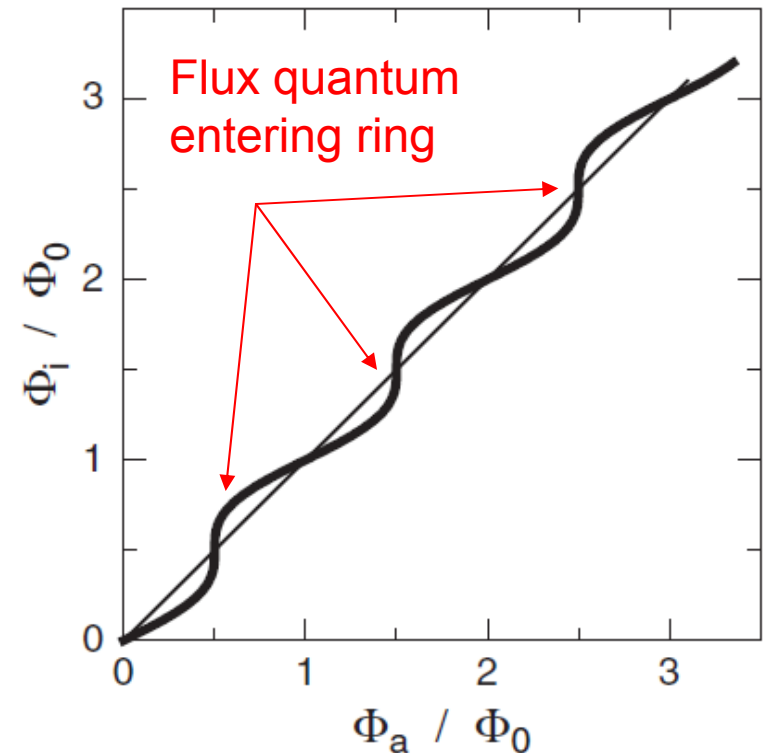
## Flux quantization (II)

- Flux quantization  $\Rightarrow$  current quantization
- Continuous variation of current is not allowed
- Phase change could be allowed, but this would temporarily destroy Cooper pair coherence, raising condensation energy
- Therefore no such jump can occur, and no flux quanta can leave the superconducting loop
- Hence persistent currents are absolutely stable

# Josephson Effect

- **Josephson junction:** thin ( $< 2\text{nm}$ ) resistor layer between two superconductors (weak link): Cooper pairs can tunnel through junction allowing current flux
- Think of ring in previous case with a weak link at one point: flux can escape from ring, one line (quantum  $\Phi_0$ ) at a time.

Graph of variation of **flux inside the ring  $\Phi_i$**  versus flux  $\Phi_a = A \bullet B$  due to external field  $B$  if ring was not superconducting, where  $A$  is the ring surface:



# Superconducting QUantum Interference Device

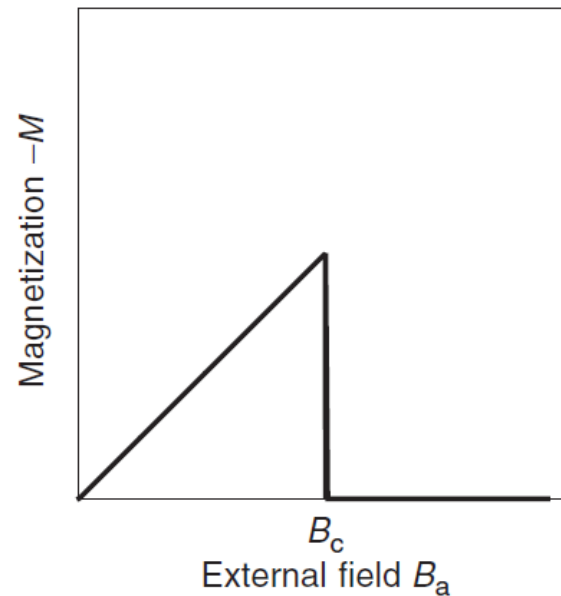
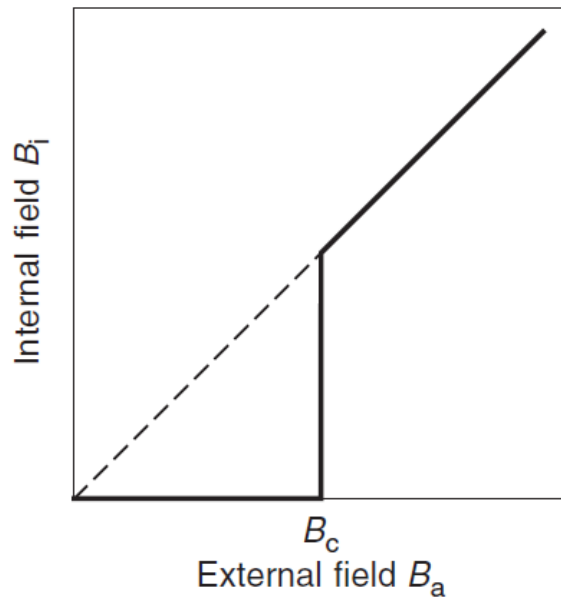
- **SQUID**: superconducting ring with one or two Josephson junctions coupled to sensitive measuring circuit
- Can measure magnetic field variations down to  $10^{-14}\text{T}$
- Most sensitive magnetic field measuring device
- Wide use since many years

Josephson effect predicted in 1962, junction first made in 1963, first squid in 1964



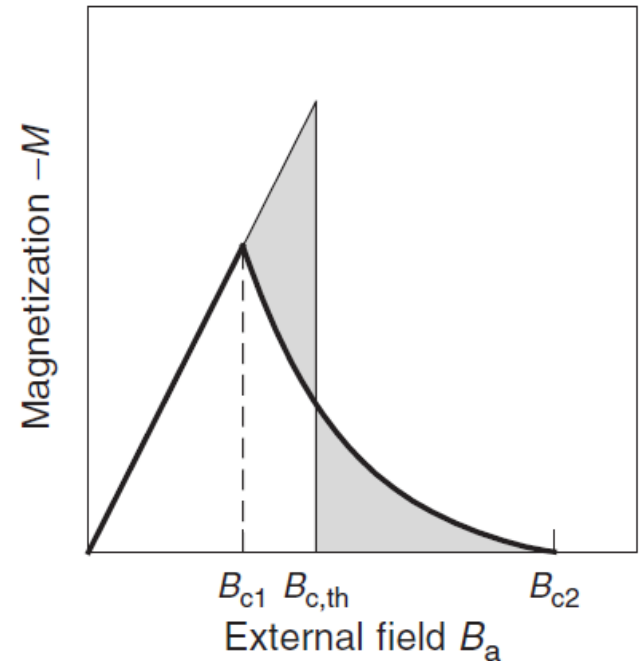
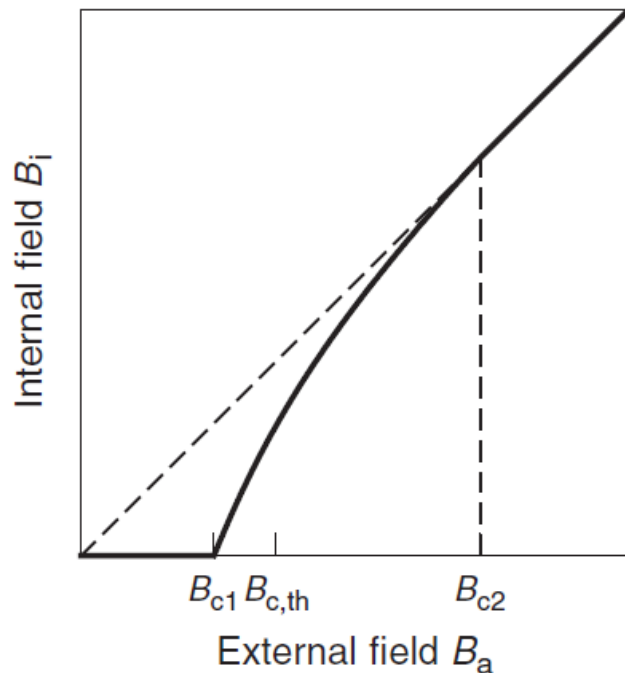
# Type I superconductors

- Most pure metals (Pb, Hg, In, Al) are type I superconductors
- They are ideal diamagnets
- In external magnetic field superconductivity breaks down above critical field  $B_c$
- See diagram of field inside superconductor  $B_i$  as function of external field  $B_a$  and the negative magnetization of the superconductor (for  $T=0$ )



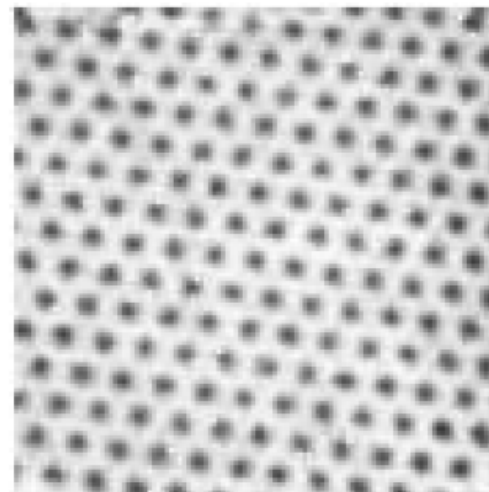
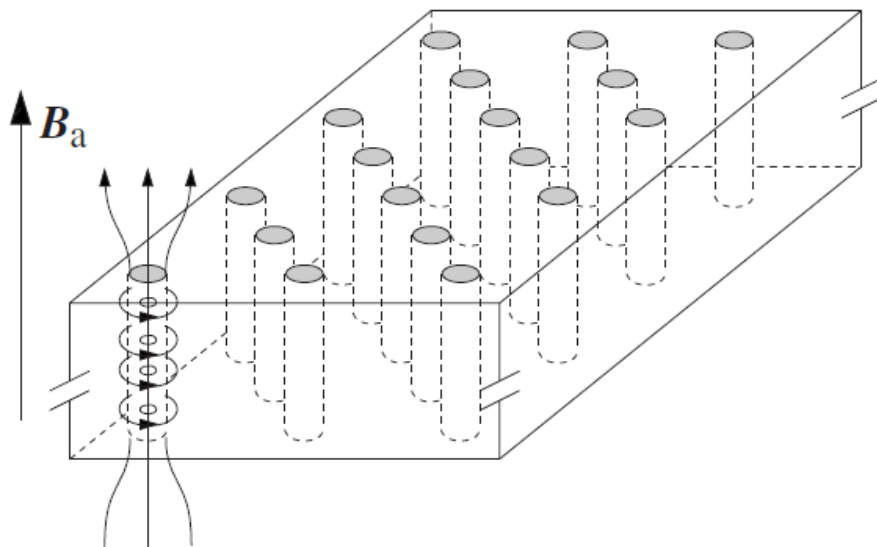
# Type II superconductors

- Transition metals, metallic glasses, novel high- $T_c$  superconductors, are type II superconductors
- Same behaviour below lower critical field:  $0 < B < B_{c1}$
- At higher field normal conducting islands are formed inside superconductive volume and field penetrates partially: mixed state for  $B_{c1} < B < B_{c2}$
- See diagram of field inside superconductor  $B_i$  as function of external field  $B_a$  and the negative magnetization of the superconductor (for  $T=0$ )



# Type II superconductors

- Type II superconductors more practical for high field applications (magnet coils): high  $B_{c2}$  and  $T_c$
- Flux lines regularly arranged in perfect crystals (Abrikosov lattice, see graph) as this minimizes energy through equal separation
- In practice sample impurities can pin flux lines



← 6000 Å →

Abrikosov lattice sketch (left) and observation (right) in  $\text{NbSe}_2$  at 1.8K and  $B=1\text{T}$ , using STM

# High $T_C$ superconductors

- For elements highest  $T_C=9.5\text{K}$  (Nb)
- For alloys highest  $T_C=23.2\text{K}$  ( $\text{Nb}_3\text{Ge}$ )

Since 1986:

- $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ :  $T_C=36\text{K}$
- $\text{YBa}_2\text{Cu}_3\text{O}_7$ :  $T_C=90\text{K}$
- $\text{Tl}_2\text{Ba}_2\text{Ca}_2\text{Cu}_3\text{O}_{10}$ :  $T_C=120\text{K}$

## High $T_C$ superconductors

- Materials difficult to make
- Cannot transport large currents
- Intriguing as mechanism is still not clear

# High $T_C$ superconductors

- Crystal structure of  $\text{CuO}_2$  layers spaced apart by other atoms
- Anisotropic behaviour - superconductivity mainly along planes
- High  $T_C$  indicates strong pairing interaction
- Small size of Cooper pairs
- April 2008:  $\text{LaOFeAs}$  (first not cuprate high  $T_C$  superconductor) with  $T_C$  up to 50K; very little theoretical understanding

